

Spectroscopic measurements of flow and ion temperature at SSX

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1. Introduction

• Ideal MHD postulates that plasma resistivity is zero, such that Ohm's Law for the plasma is

 $\mathbf{E} + \mathbf{u} \times \mathbf{B} = 0$

- Ideal MHD predicts that magnetic field lines should be "frozen in" to the plasma
- Magnetic reconnection occurs when a large field gradient results in a breakdown of ideal MHD



Figure 1: Sweet-Parker picture of 2D magnetic reconnection

- **B** consumed during reconnection; energy density $B^2/2\mu_0$ of reconnected field converted to flow and heating
- Reconnection known to be complex, 3D process
- Simplest models inadequate; Sweet-Parker normalized reconnection rate $1/\sqrt{S}$ orders of magnitude too slow
- Magnetic reconnection thought to be involved in numerous astrophysical phenomena, including the heating of the solar corona

2. The Swarthmore Spheromak Experiment

- The Swarthmore Spheromak Experiment (SSX) studies magnetic reconnection by merging co- and counter-helicity spheromaks
- Spheromaks are toroidal plasma configurations with both poloidal and toroidal magnetic fields
- Spheromaks are an example of a force-free Grad-Shafranov equilibrium, with a pressure balance between kinetic and magnetic pressure
- Typical plasma parameters in SSX include electron density $n_e \sim 10^{15} \text{ cm}^{-3}$, temperature $T_i + T_e \sim 30 \text{ eV}$, magnetic fields $|B| \sim 0.1$ T, $v_{\text{Alfvèn}} \sim 6 \text{ cm}/\mu \text{s}$, and Lundquist number $S \sim 1000$.



Figure 2: a) & b): Two views of the geometry and fields of a spheromak. c): Overview of SSX chamber

- Newly constructed high-resolution ion doppler spectroscopy (IDS) system allows better measurements of chordintegrated flows and ion temperatures
- IDS system observes spectral light emitted by impurity ions in the SSX hydrogen plasma
- The wavelength of light coming from plasma that is moving towards or away from the observer at speed v is Doppler shifted by $\Delta \lambda$, where

$$\Delta \lambda = \lambda_0 \frac{v}{c}$$

and λ_0 is the nominal wavelength of the light

• Assuming a Maxwellian velocity distribution, the lineshape is Gaussian with half-width

$$\Delta \lambda_{\rm FWHM} = \frac{2\lambda_0}{c} \sqrt{\frac{2kT_i \ln 2}{m}}$$

where m is the mass of the emitting ions

3. The IDS System

- Collection optics select 1 of 10 chords through the plasma at midplane.
- Light travels through fiber optic to 2x magnifying input optics at entrance slit of 1.33 m Czerny-Turner monochromator with 316 grooves/mm echelle grating
- Magnifying exit optics increase the size of the image at the exit plane of the spectrometer $\sim 4x$.
- 32-channel photomultiplier tube (PMT) array detector gives submicrosecond time response, but has wide (1mm) pixel width
- Wide (1mm) pixel width of PMT array is compensated for by the exit optics and by observing at high spectral order



Figure 3: Left: Photograph of SSX lab showing vacuum chamber, monochromator, and associated IDS optics. Right: Overhead view of spectrometer exit slit, magnifying optics, and PMT array. The fiber optic is visible at right.



Figure 4: Schematic of IDS system. Typical lineshape is shown at left.

4. Results and Discussion



Figure 5: 227.9 nm CIII line during counter-helicity meraing in two different shots. A time series is shown at left, while a single frame is shown at right. Velocities shown correspond to the Doppler shift of the emitting plasma.

- 227.9 nm C III line observed at 25th order, with dispersion 0.008 nm/mm
- Instrument temperature $\sim 6 \text{ eV}$
- Only Doppler broadening is significant; other effects such as pressure broadening are negligible.
- Double-peaked lineshape shown above seen consistently during counter-helicity merging, when toroidal fields reconnect
- Bi-directional flows observed on the sun by spacecraft [1]
- Previous laboratory experiments observed broadened spectral lines but interpreted them as high ion temperatures [2]

- Broadening could be due to transient flows averaged out due to insufficient time resolution
- Abel inversions to determine radial profiles of velocity and ion temperature under way
- Assuming cylindrical symmetry, given some parameter f(r), the quantity F(y) that is measured along a chord of height y from the center is given by

$$F(y) = 2\int_y^R \frac{f(r)r\,dr}{\sqrt{r^2 - y^2}}$$

• This equation can be solved for f(r):

$$f(r) = -\frac{1}{\pi} \int_{r}^{R} \frac{dF}{dy} \frac{dy}{\sqrt{y^{2} - r^{2}}}$$

• Because of derivative $\frac{dF}{du}$, this equation cannot be directly applied to noisy data



Figure 6: Abel inversion of plasma emissivity. Blue line shows inverted plasma emissivity as a function of radius. Black line shows line-integrated emissivity as a function of chord height. Red line is a smoothed cubic spline interpolation to the data.

- We have begun to do Abel inversions by interpolating F(y)with cubic splines
- Radial profiles of plasma emmissivity calculated for dipoletrapped single spheromak shots and field-reversed configurations (FRC's) formed during counter-helicity merging
- Abel inversions of other plasma parameters (v_{radial}, T_i) are under way

References

- [1] D.E. Innes et al. Bi-directional plasma jets produced by magnetic reconnection on the Sun. Nature, 386 (1993), 811.
- [2] Y. Ono et al. Experimental investigation of threecomponent magnetic reconnection by use of merging spheromaks and tokamaks. Phys. Plasmas, 4 (1997), 1953.



